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Reduced Order Modeling of a Free Boundary Osmotic Cell Swelling Problem with Exact Mass Conservation

Christoph Lehrenfeld, Stephan Rave

Outline

1. Reduced Basis Methods for Advection Dominated Problems.
2. A Globally Mass Conservative Nonlinear Reduced Basis Method for Parabolic Free Boundary Problems.

Reduced Basis Methods for Advection Dominated Problems

Parametric Model Order Reduction

Consider time-dependent parametric problems

$$\Phi : \mathcal{P} \rightarrow X([0, T]; V_h), \quad s : X([0, T]; V_h) \rightarrow \mathbb{R}^S$$

where

- ▶ $\mathcal{P} \subset \mathbb{R}^P$ parameter domain.
- ▶ V_h “truth” solution state space, $\dim V_h \gg 0$.
- ▶ Φ maps parameters to solutions (*hard* to compute).
- ▶ s maps state vectors to quantities of interest.

Objective

Compute

$$s \circ \Phi : \mathbb{R}^P \rightarrow X([0, T]; V_h) \rightarrow \mathbb{R}^S$$

for *many* $\mu \in \mathcal{P}$ or *quickly* for unknown single $\mu \in \mathcal{P}$.

Reduced Basis Methods: Three Basic Ideas

Objective

Compute

$$s \circ \Phi : \mathbb{R}^P \rightarrow X([0, T]; V_h) \rightarrow \mathbb{R}^S$$

When Φ , s sufficiently smooth, quickly computable low-dimensional approximation of $s \circ \Phi$ should exist.

▶ **Idea 1:** State space projection:

▶ Define approximation $\Phi_N : \mathcal{P} \rightarrow X([0, T]; V_N)$, $N := \dim V_N \ll \dim V_h$, via (Petrov-)Galerkin projection.

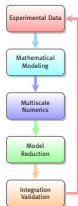
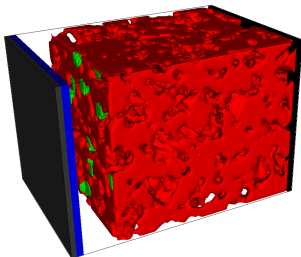
▶ Approximate $s \circ \Phi \approx s \circ \Phi_N$.

▶ **Idea 2:** Construct V_N from PODs of solution snapshots $\Phi(\mu_1), \dots, \Phi(\mu_k)$.

▶ **Idea 3:** Select μ_1, \dots, μ_k iteratively via greedy search over \mathcal{P} using quickly computable surrogate $\eta(\Phi_N(\mu), \mu) \geq \|\Phi(\mu) - \Phi_N(\mu)\|$ (POD-GREEDY).

+ **Hyper-reduction technique (EIM, DEIM, GEIM, Gappy POD, ...)**

Example: RB Approximation of Li-Ion Battery Models



MULTIBAT: Gain understanding of degradation processes in rechargeable Li-Ion Batteries through mathematical modeling and simulation at the pore scale.

FOM:

- ▶ 2.920.000 DOFs
- ▶ Simulation time: $\approx 15.5h$

ROM:

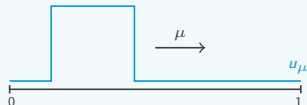
- ▶ Snapshots: 3
- ▶ $\dim V_N = 245$
- ▶ Rel. err.: $< 4.5 \cdot 10^{-3}$
- ▶ Reduction time: $\approx 14h$
- ▶ Simulation time: $\approx 8m$
- ▶ Speedup: 120

Trouble with Advection Dominated Problems

Typically slow decay of Kolmogorov N -widths d_N of the solution manifold, but RB will only work well for rapid decay!

$$d_N := \inf_{\substack{V_N \subseteq V_h \\ \dim V_N \leq N}} \sup_{\substack{u \in \Phi(\mathcal{P}) \\ t \in [0, T]}} \|u(t) - P_{V_N}(u(t))\|.$$

Basic example



$$\begin{aligned} \partial_t u(t, x) + \mu \cdot \partial_x u_\mu(t, x) &= 0 \\ u_\mu(0, x) &= u_0(x), \quad u_\mu(0, t) = u_\mu(1, t) \\ \mu, x, t &\in [0, 1] \end{aligned}$$

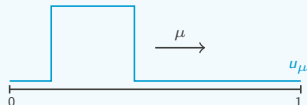
Here: $d_N \sim N^{-1/2}$ w.r.t. $L^2([0, 1])$.

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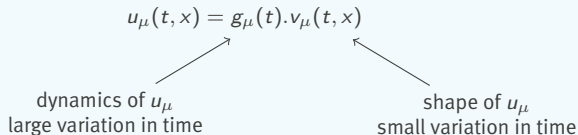
However: We can describe solution easily as

$$u_\mu(t, x) = u_0(x - \mu \cdot t \bmod 1).$$

Nonlinear Approximation

General Idea

Write $u_\mu(t, x)$ as

$$u_\mu(t, x) = g_\mu(t) \cdot v_\mu(t, x)$$


dynamics of u_μ
large variation in time

shape of u_μ
small variation in time

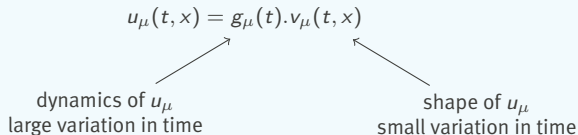
where \mathcal{V} function space, $v_\mu(t) \in \mathcal{V}$ and $g_\mu(t)$ is element of Lie group G acting on \mathcal{V} .

- ▶ $v_\mu(t, x)$ should be easier to approximate by a linear space than $u_\mu(t, x)$!

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- ▶ $v_\mu(t, x)$ should be easier to approximate by a linear space than $u_\mu(t, x)$!
- ▶ Related/other approaches: [Rowley, Marsden, 2000] [Gerbeau, Lombardi, 2014] [Iollo, Lombardi, 2014] [Carlberg, 2015] [Taddei, Perotto, Quarteroni, 2015] [Reiss, Schulze, Sesterhenn, Mehrmann, 2015] [Cagniard, Maday, Stamm, 2016] [Nair, Balajewicz, 2017] [Welper, 2017] [Rim, Moe, LeVeque, 2018] ...

Method of Freezing [Beyn, Thümmeler, 2004], [Rowley et. al., 2000, 2003]

Definition (Method of Freezing)

With initial conditions $v_\mu(0) = u(0)$, $g_\mu(0) = e$, solve:

$$\partial_t v_\mu(t) + \mathcal{L}_\mu(v_\mu(t)) + g_\mu(t) \cdot v_\mu(t) = 0$$

$$\Phi(v_\mu(t), g_\mu(t)) = 0$$

frozen PDAE

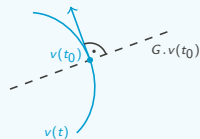
$$g_\mu(t) = g(t)_\mu^{-1} \partial_t g_\mu(t)$$

reconstruction equation

Orthogonality phase condition

$$\Phi(v, g) = 0 \iff \partial_t v(t) \perp G \cdot v(t)$$

$$\iff (\mathcal{L}(v) + g \cdot v, h \cdot v) = 0 \quad \forall h \in G$$



Test Problem

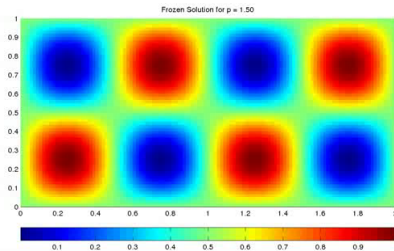
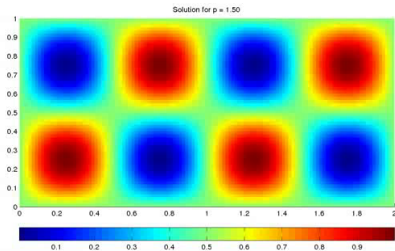
2D Burgers-type problem

Solve on $\Omega = [0, 2] \times [0, 1]$ with periodic boundaries, $t \in [0, 0.3]$, $\vec{v} \in \mathbb{R}^2$ and $\mu \in [1, 2]$:

$$\partial_t u + \nabla \cdot (\vec{v} \cdot u^\mu) = 0$$

$$u(0, x_1, x_2) = 1/2(1 + \sin(2\pi x_1) \sin(2\pi x_2))$$

Let $G := \mathbb{R}^2$ act on u by periodic shifts.



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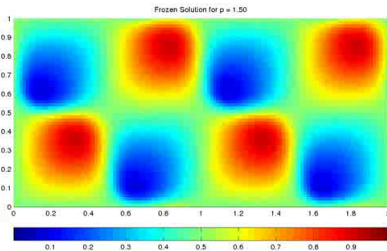
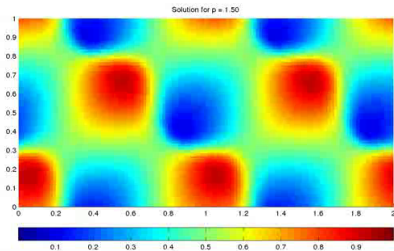
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Combining RB with the Method of Freezing

FrozenRB-Scheme for 2D-shifts [Ohlberger, R, 2013]

Solve

$$\partial_t v_{\mu(t),N} + P_{V_N} \circ \mathcal{I}_M[\mathcal{L}\mu](v_{\mu,N}(t)) - g_{\mu(t),N} \cdot (P_{V_N} \circ \nabla)(v_{\mu,N}(t)) = 0$$

$$[(\partial_{x_i} v_{\mu,N}, \partial_{x_j} v_{\mu,N})]_{i,j} \cdot [g_{\mu,N}]_j = [(\mathcal{I}_M[\mathcal{L}\mu](v_{\mu,N}), \partial_{x_i} v_{\mu,N})]_i$$

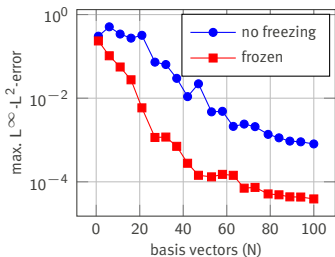
and

$$\partial_t g_{\mu}(t) = g_{\mu}(t)$$

with initial conditions $v_{\mu}(0) = u(0)$, $g_{\mu}(0) = (0, 0)^T$.

- ▶ EI-GREEDY, POD-GREEDY algorithms for basis generation.
- ▶ **Full offline/online decomposition.**
- ▶ No additional evaluations of nonlinearity.

Results for the Burgers Problem

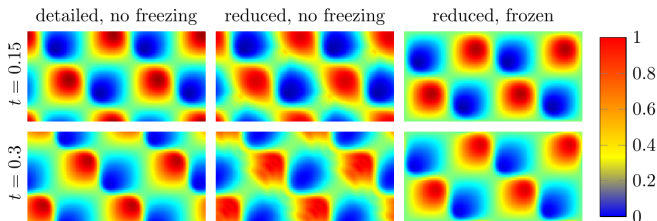


Left:

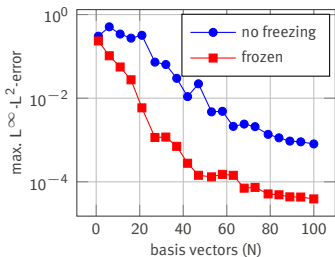
- ▶ $1.9 \cdot N$ interpolation points.
- ▶ Test set: 100 random μ .

Bottom:

- ▶ $\dim V_N = 20$, 38 interpolation points.
- ▶ $\mu = 1.5$.



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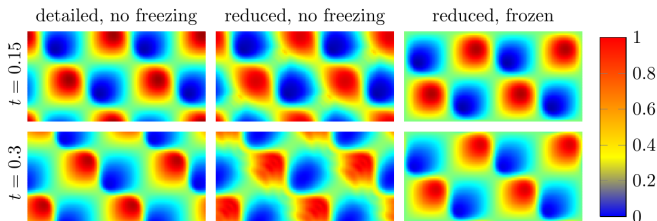


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Freezing of
multiple waves?
See Harshit
Bansal's talk at
11:05, Room 532!

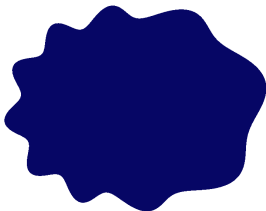
A Globally Mass Conservative Nonlinear Reduced Basis Method for Parabolic Free Boundary Problems

A Free Boundary Problem

Osmotic cell swelling model [Lippho, Prokert, 2012]

Given $\Omega(0) \subset \mathbb{R}^d$, $u(0) \in H^1(\Omega(0))$ and coefficients $u_{\text{ext}}, \alpha, \beta, \gamma \in \mathbb{R}$, the **concentration** $u(t)$ and **normal velocity** w_Γ of $\partial\Omega(t)$ is given by:

$$\begin{aligned} \partial_t u - \alpha \Delta u &= 0 && \text{in } \Omega(t), \\ w_\Gamma u + \alpha \partial_n u &= 0 && \text{on } \Gamma(t), \\ -\beta \kappa + \gamma(u - u_{\text{ext}}) &= w_\Gamma && \text{on } \Gamma(t). \end{aligned}$$



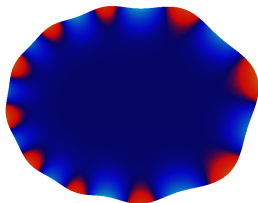
- ▶ u_0 : constant concentration in $\Omega(t)^c$
- ▶ κ : mean curvature of $\partial\Omega(t)$
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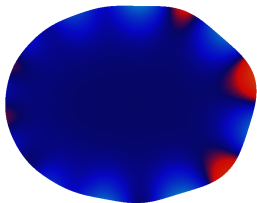
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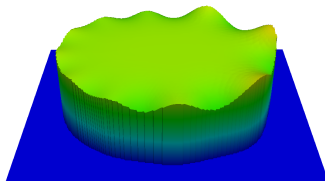
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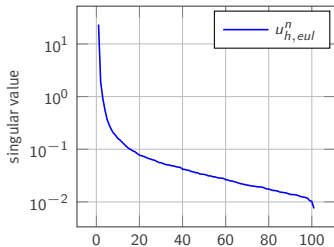
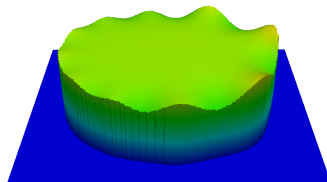
Eulerian Approximation

- ▶ Consider $u(t) \in L^2(\Omega(t)) \hookrightarrow L^2(\mathbb{R}^d)$ as joint approximation space.



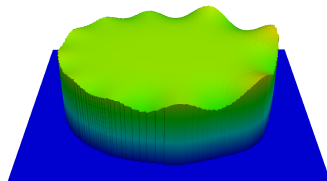
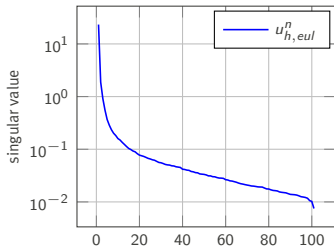
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Idea

Use nonlinear transformation

$$u(t)[\Psi(t)[x]]$$

to freeze boundary $\Gamma(t)$ in space.

- ▶ Fix reference domain

$$\hat{\Omega} := \Psi(t)^{-1}(\Omega(t)).$$

ALE Formulation

Fix reference domain $\hat{\Omega}$ and introduce deformation field $\Psi(t)$ s.t. $\Psi(t)(\hat{\Omega}) = \Omega(t)$.

Pulling back the equations to $\hat{\Omega}$ leads to the following time-discretization scheme:

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1. Compute boundary velocity:

$$\begin{aligned} & \int_{\widehat{\Gamma}} J_{\Gamma}^{n-1} \widehat{\mathbf{w}}_{\Gamma,h}^{n-1} \cdot \widehat{\mathbf{s}}_h \, d\mathbf{s} + \beta \Delta t \int_{\widehat{\Gamma}} J_{\Gamma}^{n-1} (\mathbf{P} \cdot (\mathbf{F}^{n-1})^{-T} \cdot \nabla \widehat{\mathbf{w}}_{\Gamma,h}^{n-1}) : ((\mathbf{F}^{n-1})^{-T} \nabla \widehat{\mathbf{s}}_h) \, d\mathbf{s} \\ & = -\beta \int_{\widehat{\Gamma}} J_{\Gamma}^{n-1} \mathbf{P} : (\mathbf{F}^{n-1})^{-T} \nabla_{\widehat{\Gamma}} \widehat{\mathbf{s}}_h \, d\mathbf{s} + \gamma \int_{\widehat{\Gamma}} J_{\Gamma}^{n-1} (\widehat{u}_h - u_{\text{ext}}) \widehat{\mathbf{s}}_h \cdot ((\mathbf{F}^{n-1})^{-T} \widehat{\mathbf{n}}) \, d\mathbf{s}. \end{aligned}$$

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2. Extend velocity to interior via harmonic extension:

$$-\operatorname{div}[h_{\mathcal{T}}^{-1}(\nabla \hat{\mathbf{w}}_h^{n-1} + (\nabla \hat{\mathbf{w}}_h^{n-1})^T)] = 0 \quad \text{in } \hat{\Omega}, \quad \hat{\mathbf{w}}_h^{n-1} = \hat{\mathbf{w}}_{\Gamma,h}^{n-1} \quad \text{on } \partial \hat{\Omega}.$$

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$$\Psi_h^n = \Psi_h^{n-1} + \Delta t \hat{\mathbf{w}}_h^{n-1}.$$

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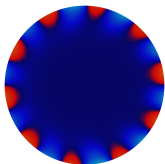
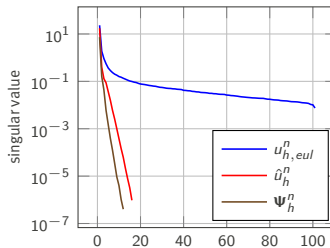
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4. Update concentration field:

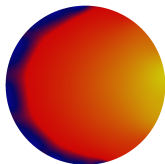
$$\begin{aligned} & \int_{\hat{\Omega}} J^n \hat{u}_h^n \hat{v}_h \, dx + \Delta t \int_{\hat{\Omega}} J^n \hat{u}_h^n \hat{\mathbf{w}}_h^{n-1} \cdot ((\mathbf{F}^n)^{-T} \cdot \nabla \hat{v}_h) \, dx \\ & + \alpha \Delta t \int_{\hat{\Omega}} J^n ((\mathbf{F}^n)^{-T} \nabla \hat{u}_h^n) \cdot ((\mathbf{F}^n)^{-T} \nabla \hat{v}_h) \, dx = \int_{\hat{\Omega}} J^{n-1} \hat{u}_h^{n-1} \hat{v}_h \, dx. \end{aligned}$$

ALE Formulation

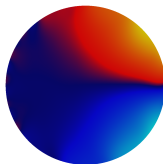
- ▶ Rapid singular value decay of both concentration and deformation fields.
- ▶ After space discretization this corresponds to moving-mesh approach (ALE), where $\Psi_h^n(v)$ is the trajectory of the vertex v .
- ▶ In contrast to “parameterized domain problems”, the domain deformation Ψ_h^n is part of the equation system.



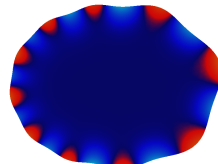
\hat{u}_h^n



Ψ_{h1}^n



Ψ_{h2}^n



reconstruction

Nonlinear RBM for Free Boundary Problems

Use standard RB machinery to construct ROM:

- ▶ Compute low-rank approximation spaces for \hat{u}_h^n , Ψ_h^n , $\hat{w}_{\Gamma,h}^n$ via POD.
(Could also use POD-GREEDY).
- ▶ Use EIM to approximate coefficient functions, vectors, tensors depending nonlinearly on Ψ_h^n .
- ▶ Similar to [Ballarin, Rozza, 2016] in context of FSI.

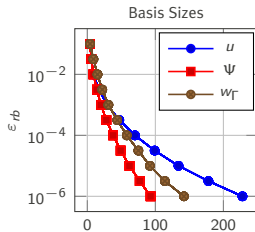
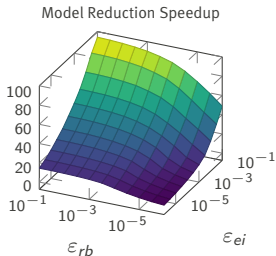
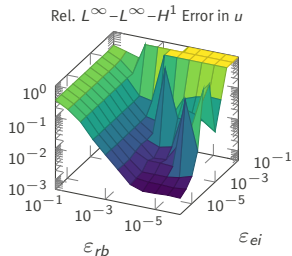
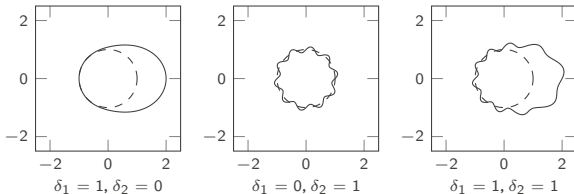
Numerical Experiment

► Parameterization:

- $\alpha \in [0.1, 1]$
- $\beta \in [0.001, 0.1]$
- $\delta_1, \delta_2 \in [0, 1]$

► Snapshots: 3^4

► FOM: 3988 + 7976 DOFS



Global Mass Conservation

Concentration update

$$\int_{\hat{\Omega}} J^n \hat{u}_h^n \hat{v}_h \, d\mathbf{x} + \Delta t \int_{\hat{\Omega}} J^n \hat{u}_h^n \hat{\mathbf{w}}_h^{n-1} \cdot ((\mathbf{F}^n)^{-T} \cdot \nabla \hat{v}_h) \, d\mathbf{x} \\ + \alpha \Delta t \int_{\hat{\Omega}} J^n ((\mathbf{F}^n)^{-T} \nabla \hat{u}_h^n) \cdot ((\mathbf{F}^n)^{-T} \nabla \hat{v}_h) \, d\mathbf{x} = \int_{\hat{\Omega}} J^{n-1} \hat{u}_h^{n-1} \hat{v}_h \, d\mathbf{x}.$$

Global Mass Conservation

Concentration update

$$\begin{aligned}
 \int_{\hat{\Omega}} J^n \hat{u}_h^n \hat{v}_h \, d\mathbf{x} + \Delta t \int_{\hat{\Omega}} J^n \hat{u}_h^n \hat{\mathbf{w}}_h^{n-1} \cdot ((\mathbf{F}^n)^{-T} \cdot \nabla \hat{v}_h) \, d\mathbf{x} \\
 + \alpha \Delta t \int_{\hat{\Omega}} J^n ((\mathbf{F}^n)^{-T} \nabla \hat{u}_h^n) \cdot ((\mathbf{F}^n)^{-T} \nabla \hat{v}_h) \, d\mathbf{x} = \int_{\hat{\Omega}} J^{n-1} \hat{u}_h^{n-1} \hat{v}_h \, d\mathbf{x}.
 \end{aligned}$$

► Testing with $\hat{v}_h \equiv 1$ yields:

$$\int_{\Omega^n} u_h^n \, d\mathbf{x} = \int_{\hat{\Omega}} J^n \hat{u}_h^n \, d\mathbf{x} + 0 + 0 = \int_{\hat{\Omega}} J^{n-1} \hat{u}_h^n = \int_{\Omega^{n-1}} u_h^{n-1} \, d\mathbf{x}$$

Global Mass Conservation

Concentration update

$$\int_{\hat{\Omega}} J^n \hat{u}_h^n \hat{v}_h \, dx + \Delta t \int_{\hat{\Omega}} J^n \hat{u}_h^n \hat{w}_h^{n-1} \cdot ((\mathbf{F}^n)^{-T} \cdot \nabla \hat{v}_h) \, dx$$

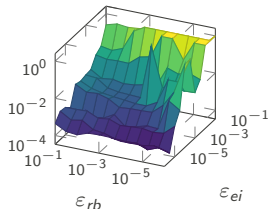
$$+ \alpha \Delta t \int_{\hat{\Omega}} J^n ((\mathbf{F}^n)^{-T} \nabla \hat{u}_h^n) \cdot ((\mathbf{F}^n)^{-T} \nabla \hat{v}_h) \, dx = \int_{\hat{\Omega}} J^{n-1} \hat{u}_h^{n-1} \hat{v}_h \, dx.$$

- ▶ Testing with $\hat{v}_h \equiv 1$ yields:

$$\int_{\Omega^n} u_h^n \, dx = \int_{\hat{\Omega}} J^n \hat{u}_h^n \, dx + 0 + 0 = \int_{\hat{\Omega}} J^{n-1} \hat{u}_h^n = \int_{\Omega^{n-1}} u_h^{n-1} \, dx$$

- ▶ Mass conservation is preserved by RB projection by adding 1 to RB for u_h^n .
- ▶ Inexact assembly of mass matrix due to EI destroys mass conservation.

Rel. Mass Conservation Error



Global Mass Conservation with EI

- ▶ Note that in 2D:

$$\int_{\hat{\Omega}} J^n \hat{u}_h^n \hat{v}_h = m(\Psi_h^n, \Psi_h^n, \hat{u}_h^n, \hat{v}_h),$$

where

$$m(\Phi_h^n, \Psi_h^n, \hat{u}_h^n, \hat{v}_h) = \int_{\hat{\Omega}} \partial_x \Phi_{hx}^n \cdot \partial_y \Psi_{hy}^n \cdot \hat{u}_h^n \cdot \hat{v}_h + \partial_x \Phi_{hy}^n \cdot \partial_y \Psi_{hx}^n \cdot \hat{u}_h^n \cdot \hat{v}_h dx.$$

Global Mass Conservation with EI

- ▶ Note that in 2D:

$$\int_{\hat{\Omega}} J^n \hat{u}_h^n \hat{v}_h = m(\Psi_h^n, \Psi_h^n, \hat{u}_h^n, \hat{v}_h),$$

where

$$m(\Phi_h^n, \Psi_h^n, \hat{u}_h^n, \hat{v}_h) = \int_{\hat{\Omega}} \partial_x \Phi_{hx}^n \cdot \partial_y \Psi_{hy}^n \cdot \hat{u}_h^n \cdot \hat{v}_h + \partial_x \Phi_{hy}^n \cdot \partial_y \Psi_{hx}^n \cdot \hat{u}_h^n \cdot \hat{v}_h dx.$$

- ▶ Could assemble mass matrix 4-tensor exactly.
- ▶ Relatively expensive.
(dim RB = 30 \implies 6MB for reduced tensor)
- ▶ 5-tensor in 3D!

Global Mass Conservation with EI

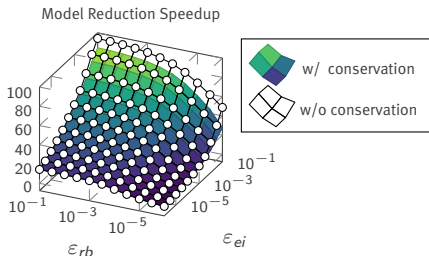
- ▶ Note that in 2D:

$$\int_{\hat{\Omega}} J^n \hat{u}_h^n \hat{v}_h = m(\Psi_h^n, \Psi_h^n, \hat{u}_h^n, \hat{v}_h),$$

where

$$m(\Phi_h^n, \Psi_h^n, \hat{u}_h^n, \hat{v}_h) = \int_{\hat{\Omega}} \partial_x \Phi_{hx}^n \cdot \partial_y \Psi_{hy}^n \cdot \hat{u}_h^n \cdot \hat{v}_h + \partial_x \Phi_{hy}^n \cdot \partial_y \Psi_{hx}^n \cdot \hat{u}_h^n \cdot \hat{v}_h dx.$$

- ▶ Could assemble mass matrix 4-tensor exactly.
- ▶ Relatively expensive.
(dim RB = 30 \implies 6MB for reduced tensor)
- ▶ 5-tensor in 3D!
- ▶ Better approach:
 1. Assemble mass matrix using EI.
 2. Assemble 3-tensor $m(\Phi_h^n, \Psi_h^n, \hat{u}_h^n, 1)$ exactly and set corresponding row of mass matrix.



Thank you for your attention!

Lehrenfeld, R, *Mass conservative reduced order modeling of a free boundary osmotic cell swelling problem*, Adv Comput Math (2019).

Ohlberger, R, *Nonlinear reduced basis approximation of parameterized evolution equations via the method of freezing*, C. R. Math. Acad. Sci. Paris, 351 (2013).

Ohlberger, R, *Reduced Basis Methods: Success, Limitations and Future Challenges*, Proceedings of ALGORITHMY 2016.

Milk, R, Schindler, *pyMOR – Generic Algorithms and Interfaces for Model Order Reduction* SIAM J. Sci. Comput., 38(5), 2016.

<http://www.pymor.org/>

My homepage (with FrozenRB code)

<http://stephanrave.de/>

Outlook: Remeshing

Strongly anisotropic mesh deformations in ALE schemes lead to:

- ▶ bad approximation spaces.
- ▶ ill-conditioned system matrices.

Possible MOR approach:

- ▶ In FOM: Locally adapt mesh $\hat{\mathcal{T}}_h$ on $\hat{\Omega}$ s.t. $\Psi_h^n(\hat{\mathcal{T}}_h)$ has good shape regularity properties.
- ▶ Solve extension problem for $\hat{\mathbf{w}}_{\Gamma,h}^n$ on Ω^n instead of $\hat{\Omega}$.
- ▶ Use “RB for AFEM” methods to construct ROM [Ullmann, Rotkvic, Lang, 2016] [Yano 2016] [Ali, Steih, Urban, 2017] [Hinze, Gräßle, 2017].
- ▶ Deformation-dependent norms?
- ▶ Dictionary-based approaches?