

Multiscale models for glioma invasion

Christina Surulescu

Glioma invasion: Microscale, mesoscale, macroscale

More genera models and theory issues

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More genera models and theory issues Joint work with:

- Christian Engwer & Markus Knappitsch (WWU Münster)
- Thomas Lorenz (UAS Rhein-Main Wiesbaden)
- Alex Hunt (TU Kaiserslautern)

Data and biomedical informations:

- Carsten Wolters & Felix Lucka (IBB, WWU Münster)
- Katarina Wolf (Radboud Univ. Nijmegen Medical Centre)
- Yvonne Dzierma (Uniklinikum des Saarlandes)



Glioblastoma multiforme (GBM)

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Diffusion tensor imaging

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Proliferation via cell-tissue interactions

Alternative proliferation modeling: Go-or-grow Therapy

More general models and theory issues

- Variant of diffusion-weighted magnetic resonance imaging, DW-MRI.
- Measures the spatial diffusion of water molecules by MRI per volume element (voxel).



MRI-device (Philips Chieva 3.0 T) http://upload.wikimedia.org/

Consider a single voxel:



This leads to a diffusion tensor

$$\mathbb{D}(\mathbf{x}) = \begin{pmatrix} d_{\scriptscriptstyle XX}(\mathbf{x}) & d_{\scriptscriptstyle XY}(\mathbf{x}) & d_{\scriptscriptstyle XZ}(\mathbf{x}) \\ d_{\scriptscriptstyle YX}(\mathbf{x}) & d_{\scriptscriptstyle YY}(\mathbf{x}) & d_{\scriptscriptstyle YZ}(\mathbf{x}) \\ d_{\scriptscriptstyle ZX}(\mathbf{x}) & d_{\scriptscriptstyle ZY}(\mathbf{x}) & d_{\scriptscriptstyle ZZ}(\mathbf{x}) \end{pmatrix}$$



Representation of anisotropic diffusion tensor data (see Hagmann et al. 2006)

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More general models and theory issues Assume diffusion tensor in diagonal form: $\mathbb{D}(\mathbf{x}) = \begin{pmatrix} \lambda_1 & 0 & 0\\ 0 & \lambda_2 & 0\\ 0 & 0 & \lambda_3 \end{pmatrix}$.























Representation of the anisotropic DTI data

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More general models and theory issues Glyphs & tractography





Fractional anisotropy



$$FA(\mathbf{x}) = \sqrt{\frac{3}{2}} \frac{\sqrt{\sum_{i=1}^{3} (\lambda_i - \bar{\lambda})^2}}{\sqrt{\sum_{i=1}^{3} \lambda_i^2}}$$



Modeling scales



Goal: multiscale descriptions



Biochemical basis of the microscopic model

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More general models and theory issues • Cells interact with the neighbouring tissue in order to move forward (contact guidance)



• Receptor binding to unsoluble components Q

Notation:
$$y := RQ$$
.



Biochemical basis of the microscopic model

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Notation:

• Cells interact with the neighbouring tissue in order to move forward (contact guidance)



• Receptor binding to unsoluble components Q

$$Q + (R_0 - y) \stackrel{k^+}{\underset{k^-}{\rightleftharpoons}} RQ$$
$$y := RQ.$$



Glioma follow white matter tracts (Giese & Westphal, 1996)

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Glia cells and myelinated axon bundles



TRENDS in Neurosciences



Individual variables (N cells)

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Position	$\mathbf{x}^{(j)} \in \mathbb{R}^n$
Velocity	$\mathbf{v}^{(j)} \in V = s \mathbb{S}^{n-1}$
Receptor	state $y^{(j)} \in Y \ (j = 1,, N).$

• Newton's law (in the absence of reorientations)

$$rac{d\mathbf{x}^{(j)}}{dt}=\mathbf{v}^{(j)}$$
 , $rac{d\mathbf{v}^{(j)}}{dt}=0$

• ODE for receptor dynamics

$$\frac{dy^{(j)}}{dt} = G(y^{(j)}, Q(t, \mathbf{x}^{(j)}))$$



Micro-meso model: kinetic transport equations

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More general models and theory issues Wanted: cell density $p(t, \mathbf{x}, \mathbf{v}, y)$ at time t, position $\mathbf{x} \in \mathbb{R}^n$, with velocity $\mathbf{v} \in V \subset \mathbb{R}^n$, and internal state $y \in Y \subseteq [0, R_0]$.

$$\partial_t p + \nabla_{\mathbf{x}} \cdot (\mathbf{v}p) + \partial_y \left(G(y)p \right) = -\lambda(y)p + \lambda(y) \int_V K(\mathbf{x}, \mathbf{v}, \mathbf{v}')p(\mathbf{v}')d\mathbf{v}'$$

We choose $K(\mathbf{x}, \mathbf{v}, \mathbf{v}') = \frac{q(\mathbf{x}, \hat{\mathbf{v}})}{\omega}$, where $q(\mathbf{x}, \theta)$ represents the directional distribution of tissue fibers and $\omega = \int_V q(\hat{\mathbf{v}}) d\mathbf{v} = s^{n-1}$ where $V = s \mathbb{S}^{n-1}$.

• $q(\mathbf{x},\cdot)\in L^2(\mathbb{S}^{n-1}),\;q(\mathbf{x}, heta)\geq 0$,

• For unoriented tissue $\mathbb{E}_q(\mathbf{x}) = \int_{\mathbb{S}^{n-1}} \theta q(\mathbf{x}, \theta) d\theta = \mathbf{0}$ Notation for turning operator

$$\mathcal{L}[\lambda(y)] p := -\lambda(y) p + \lambda(y) rac{q(\mathbf{x}, \hat{\mathbf{v}})}{\omega} \int_V p(\mathbf{v}') d\mathbf{v}'.$$



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More general models and theory issues Subcellular dynamics: $\frac{d}{dt}y(t) = G(y(t), Q(t, \mathbf{x})).$

 $Q(t, \mathbf{x})$: volume fraction of tissue fibres.

$$\dot{y}=k^+(R_0-y)Q-k^-y.$$

Together with the kinetic transport equation

$$\partial_t p + \nabla_{\mathbf{x}} \cdot (\mathbf{v}p) + \partial_y (G(y)p) = \mathcal{L}[\lambda(y)]p$$

this leads to a micro-meso model (the tissue informations Q and q serve as input).

Issues:

- Proliferation not included.
- Numerical handling too complicated.
- Actually wanted: macroscopic cell density.



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Steady state:
$$y^* = \frac{k^+ Q R_0}{k^+ Q + k^-}$$
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Steady state:
$$y^* = \frac{k^+ Q R_0}{k^+ Q + k^-}$$
.

Introduce a new internal variable $z := y^* - y$ measuring deviations from the steady state.



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Consider the path of a single cell starting in \mathbf{x}_0 and moving with velocity \mathbf{v} through a time-invariant density field $Q(\mathbf{x})$.



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Turning rate: $\lambda(z) = \lambda_0 - \lambda_1 z \ge 0$, with adequate λ_0 , $\lambda_1 > 0$.



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Turning rate: $\lambda(z) = \lambda_0 - \lambda_1 z \ge 0$, with adequate λ_0 , $\lambda_1 > 0$.

Wanted: macroscopic cell density

$$M(t,\mathbf{x}) := \iint_{V \times Z} p(t,\mathbf{x},\mathbf{v},z) dz d\mathbf{v},$$

where $Z \subseteq [R_0 - y^*, y^*]$ is the shifted domain for internal states.



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More general models and theory issues Notations: $\mathbf{x} = \mathbf{x}_0 + \mathbf{v}t$, $f(Q(\mathbf{x})) = \frac{k^+Q(\mathbf{x})R_0}{k^+Q(\mathbf{x})+k^-}$ Then for any t

$$\frac{d}{dt}f(Q(\mathbf{x}_0+\mathbf{v}t))=f'(Q(\mathbf{x}_0+\mathbf{v}t))\mathbf{v}\cdot\nabla Q(\mathbf{x}_0+\mathbf{v}t)$$

and hence

$$\begin{aligned} \dot{z} &= -(k^+Q(\mathbf{x})+k^-)z + f'(Q(\mathbf{x}))\mathbf{v}\cdot\nabla Q(\mathbf{x}) \\ &= -(k^+Q(\mathbf{x})+k^-)z + \frac{k^+k^-R_0}{(k^+Q(\mathbf{x})+k^-)^2} \,\mathbf{v}\cdot\nabla Q(\mathbf{x}). \end{aligned}$$

Assume for simplicity that the receptor dynamics equilibrates rapidly, s.t. the system is close to steady-state. Then

$$\partial_t p + \mathbf{v} \cdot \nabla p - \partial_z (((k^+ Q + k^-)z - f'(Q)\mathbf{v} \cdot \nabla Q)p) = \mathcal{L}[\lambda_0]p + \mathcal{L}[\lambda_1]zp$$



Modeling proliferation: interactions with tissue

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$$\mathcal{P}(p) = \underbrace{\mu(\mathbf{x}, M, \mathbf{v})}_{\text{growth rate}} \int_{Z} \chi(\mathbf{x}, z, z') p(t, \mathbf{x}, \mathbf{v}, z') Q(\mathbf{x}) dz',$$

with $M(t, \mathbf{x}) = \int_{V} \int_{Z} p(t, \mathbf{x}, \mathbf{v}, z) dz d\mathbf{v}.$

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h $M(t, \mathbf{x}) = \int_{V} \int_{Z} p(t, \mathbf{x}, \mathbf{v}, z) dz d\mathbf{v}.$

The kernel χ characterizes the transition from the state z' to the state z during a proliferative action (cell-tissue interaction).

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The kernel χ characterizes the transition from the state z' to the state z during a proliferative action (cell-tissue interaction). Then our kinetic transport equation (KTE) becomes

$$\partial_t p + \nabla_{\mathbf{x}} \cdot (\mathbf{v}p) - \partial_z (((k^+Q + k^-)z - f'(Q)\mathbf{v} \cdot \nabla Q)p) \\ = \mathcal{L}[\lambda_0]p + \mathcal{L}[\lambda_1]zp + \mathcal{P}(p).$$

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$$\mathcal{P}(p) = \underbrace{\mu(\mathbf{x}, M, \mathbf{v})}_{\text{growth rate}} \int_{Z} \chi(\mathbf{x}, z, z') p(t, \mathbf{x}, \mathbf{v}, z') Q(\mathbf{x}) dz',$$

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$$\partial_t p + \nabla_{\mathbf{x}} \cdot (\mathbf{v}p) - \partial_z (((k^+Q + k^-)z - f'(Q)\mathbf{v} \cdot \nabla Q)p) \\ = \mathcal{L}[\lambda_0]p + \mathcal{L}[\lambda_1]zp + \mathcal{P}(p).$$

Remark: KTAP by N. Bellomo assumes cell-cell interactions

$$P_i[p](t,y) = \sum_{h,k=1}^n \mu_{hk} \iint_{Y \times Y} \chi^i_{hk}(y',y'';y) p_h(t,y') p_k(t,y'') dy' dy''.$$



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Moments

$$m(t, \mathbf{x}, \mathbf{v}) = \int_{Z} p(t, \mathbf{x}, \mathbf{v}, z) dz,$$
$$M(t, \mathbf{x}) = \iint_{V \times Z} p(t, \mathbf{x}, \mathbf{v}, z) dz d\mathbf{v}$$
$$m^{z}(t, \mathbf{x}, \mathbf{v}) = \int_{Z} zp(t, \mathbf{x}, \mathbf{v}, z) dz,$$
$$M^{z}(t, \mathbf{x}) = \iint_{V \times Z} zp(t, \mathbf{x}, \mathbf{v}, z) dz d\mathbf{v}$$

Higher order moments are set to zero, due to the small deviations z from the steady-state y^* .

We assume the data to be compactly supported in the $(\mathbf{x}, \mathbf{v}, z)$ space; this allows to perform the computations for the subsequent macroscopic scaling.



Macroscopic scaling

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More general models and theory issues • Integrate the KTE w.r.t. z to obtain

$$\partial_t m + \nabla_{\mathbf{x}} \cdot (\mathbf{v}m) = -\lambda_0 m + \lambda_1 m^z + \lambda_0 \frac{q}{\omega} M - \lambda_1 \frac{q}{\omega} M^z + \mu(\mathbf{x}, M, \mathbf{v}) Q(\mathbf{x}) \int_Z \int_Z \chi(\mathbf{x}, z, z') p(z') dz' dz$$

• Multiply the KTE by z and integrate w.r.t. z to obtain

$$\partial_t m^z + \nabla_{\mathbf{x}} \cdot (\mathbf{v} m^z) = -(k^+ Q + k^- + \lambda_0) m^z + f'(Q) \mathbf{v} \cdot \nabla Q m + \lambda_0 \frac{q}{\omega} M^z + \mu(\mathbf{x}, M, \mathbf{v}) Q(\mathbf{x}) \int_Z \int_Z z \chi(\mathbf{x}, z, z') p(z') dz' dz.$$



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More genera models and theory issues **Parabolic scaling:** $\hat{t} = \varepsilon^2 t$, $\hat{\mathbf{x}} = \varepsilon \mathbf{x}$ leads to (drop the hats!):

$$\begin{aligned} \epsilon^{2}\partial_{t}m + \epsilon\nabla\cdot(\mathbf{v}m) &= -\lambda_{0}m + \lambda_{1}m^{z} + \lambda_{0}\frac{q}{\omega}M - \lambda_{1}\frac{q}{\omega}M^{z} \\ &+ \epsilon^{2}\mu(\mathbf{x}, M, \mathbf{v})Q(\mathbf{x})\int_{Z}\int_{Z}\chi(\mathbf{x}, z, z')p(z') dz' dz \\ \epsilon^{2}\partial_{t}m^{z} + \epsilon\nabla\cdot(\mathbf{v}m^{z}) &= -(k^{+}Q + k^{-} + \lambda_{0})m^{z} + \lambda_{0}\frac{q}{\omega}M^{z} \\ &+ \epsilon f'(Q)\mathbf{v}\cdot\nabla Qm \\ &+ \epsilon^{2}\mu(\mathbf{x}, M, \mathbf{v})Q(\mathbf{x})\int_{Z}\int_{Z}Z\chi(\mathbf{x}, z, z')p(z') dz' dz. \end{aligned}$$

Remark. The proliferation rate is rescaled with ϵ^2 to let it act on the correct new time scale.



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More general models and theory issues Hilbert expansions of the moments:

$$m = \sum_{k=0}^{\infty} \epsilon^{k} m_{k} \qquad \qquad M = \sum_{k=0}^{\infty} \epsilon^{k} M_{k}$$
$$m^{z} = \sum_{k=0}^{\infty} \epsilon^{k} m_{k}^{z} \qquad \qquad M^{z} = \sum_{k=0}^{\infty} \epsilon^{k} M_{k}^{z}$$

Collect the coefficients of the powers of ϵ and integrate w.r.t. ${\bf v}$ to obtain:

$$\begin{split} m_0 &= \frac{q}{\omega} M_0, \qquad m_0^z = 0 \\ M_0^z &= 0, \qquad \qquad m_1 = -\frac{1}{\lambda_0} \left(\nabla \cdot \left(\mathbf{v} \frac{q}{\omega} M_0 \right) - \lambda_1 m_1^z \right) \\ m_1^z &= \frac{f'(Q)}{k^+ Q + k^- + \lambda_0} \nabla \cdot \left(\mathbf{v} Q \frac{q}{\omega} \right) M_0, \\ M_1 &= 0, \qquad \qquad M_1^z = 0. \end{split}$$



Effective equations on macroscale

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$$\partial_t M_0 -
abla \cdot (\mathbb{D}_T(\mathbf{x}) \nabla M_0) + \nabla \cdot (g(Q(\mathbf{x})) \mathbb{D}_T(\mathbf{x}) \nabla Q(\mathbf{x}) - \mathbf{u}(\mathbf{x})) M_0)$$

= $Q(\mathbf{x}) \mu(\mathbf{x}, M_0) M_0,$

with

$$g(Q(\mathbf{x})) = \lambda_1 (k^+ Q + k^- + \lambda_0)^{-1} f'(Q(\mathbf{x})), \text{ where}$$

$$f(Q(\mathbf{x})) = \frac{k^+ Q(\mathbf{x}) R_0}{k^+ Q(\mathbf{x}) + k^-} \text{ (subcellular level information),}$$

$$\mathbf{u}(\mathbf{x}) = \frac{1}{\lambda_0 \omega} \int_V \mathbf{v} \otimes \mathbf{v} \nabla q \, d\mathbf{v} \quad \text{(drift velocity)}$$

$$\mathbb{D}_T(\mathbf{x}) = \frac{1}{\lambda_0 \omega} \int_V q \mathbf{v} \otimes \mathbf{v} \, d\mathbf{v} \quad \text{(tumor diffusion tensor)}$$

Semilinear advection-diffusion equation \rightsquigarrow globally well posed (with Neumann BCs and adequate ICs).



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Choices of $Q(\mathbf{x})$:



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• Fractional anisotropy (FA), from data.



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Choices of $Q(\mathbf{x})$:

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Problem: assumes Q to be high where tissue is strongly aligned \sim also true in regions of isotropic (non-aligned) and densely packed tissue??



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• Estimated Q via free path length from diffusivity measured by DTI:



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• Estimated Q via free path length from diffusivity measured by DTI:

Characteristic (diffusion) length:

$$l_c = \sqrt{Dt_c},$$

with D a diffusion-related coefficient and t_c the characteristic (diffusion) time.



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• Estimated *Q* via free path length from diffusivity measured by DTI:

Characteristic (diffusion) length:

$$l_c = \sqrt{Dt_c},$$

with D a diffusion-related coefficient and t_c the characteristic (diffusion) time. Choice of D: $tr(\mathbb{D}_W)$.



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This estimate is valid for $\mathcal{N}(0, t - s)$ -distributed increments, and ours are $\mathcal{N}(0, \sigma \cdot (t - s))$ -distributed, where σ is some estimation of the diffusion speed. We choose $\sigma = l_1$, where l_1 is the largest eigenvalue of \mathbb{D}_W .


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Volume fraction of tissue fibers: Characteristic length $l_c = \sqrt{\frac{h^2 tr(\mathbb{D}_W)}{4l_1}}$.



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Volume fraction of tissue fibers: Characteristic length $l_c = \sqrt{\frac{h^2 tr(\mathbb{D}_W)}{4l_1}}$. The free volume fraction of one voxel is l_c^3/h^3 . So the occupied volume is $Q = 1 - \frac{l_c^3}{h^3}$.



Volume fraction of tissue fibers

Multiscale models for glioma invasion

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Glioma invasion: Microscale, mesoscale, macroscale

Proliferation via cell-tissue interactions

Alternative proliferation modeling: Go-or-grow Therapy

More genera models and theory issues





Estimated Q



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More general models and theory issues

Choices of $q(\mathbf{x}, \theta)$ ($\theta \in \mathbb{S}^{n-1}$ gives the fiber orientation):



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More general models and theory issues Choices of $q(\mathbf{x}, \theta)$ ($\theta \in \mathbb{S}^{n-1}$ gives the fiber orientation):

- peanut (Hillen): $q(\mathbf{x}, \theta) = \frac{n}{|\mathbb{S}^{n-1}| \operatorname{tr} \mathbb{D}_{W}(\mathbf{x})} \theta^{t} \mathbb{D}_{W}(\mathbf{x}) \theta$. **Advantage:** very simple, convenient for calculations. **Drawback:** cannot resolve crossings fiber tracts
- bimodal von Mises-Fisher (Painter & Hillen, 2013):

 $q(\mathbf{x}, \theta) = \frac{k(\mathbf{x})}{8\pi \sinh(k(\mathbf{x}))} \left(\exp(k(\mathbf{x})\phi(\mathbf{x}) \cdot \theta) + \exp(-k(\mathbf{x})\phi(\mathbf{x}) \cdot \theta) \right)$ with $k(\mathbf{x}) = \kappa FA(\mathbf{x})$ and ϕ the leading eigenvector of the water diffusion tensor for each voxel. **Drawbacks:**

- concentration parameter κ cannot be determined by measurements (ought to be assessed from many different DTI data sets)
- FA is not satisfactory enough as indicator for anisotropy



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More general models and theory issues • orientation distribution function (ODF):

$$q(\mathbf{x}, \theta) = ODF(\theta) := \int_{0}^{\infty} \Pi(r\theta) r^2 dr,$$

describes the probability of diffusion in direction θ , where $\Pi(r\theta)$: displacement probability of a spatial point in spherical coordinates.

It can be shown (Aganj et al 2010) that

$$q(\mathbf{x}, \theta) = \frac{1}{4\pi |\mathbb{D}_W(\mathbf{x})| \left(\theta^t \mathbb{D}_W(\mathbf{x})^{-1} \theta\right)^{\frac{3}{2}}}$$

Advantages:

- available for different medical imaging techniques, including Q-Ball and HARDI;
- allows to use medical data in different forms;
- does not need supplementary parameters which are difficult to assess.



Example: peanut distribution

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$$q(\mathbf{x}, \theta) = rac{n}{|\mathbb{S}^{n-1}| \operatorname{tr} \mathbb{D}_W(\mathbf{x})} \theta^t \mathbb{D}_W(\mathbf{x}) \theta.$$

Tumor diffusion tensor:

$$\mathbb{D}_{T}(\mathbf{x}) = \frac{s^{2}}{\omega\lambda_{0}(n+2)} \left(\mathbb{I}_{n} + 2\frac{\mathbb{D}_{W}}{\operatorname{tr} \mathbb{D}_{W}} \right),$$

Drift velocity:

$$\mathbf{u}(\mathbf{x}) = \frac{s^2}{\omega\lambda_0(n+2)} \bigg[-\frac{1}{(\operatorname{tr}\,\mathbb{D}_W)^2} \bigg(\operatorname{tr}\,\mathbb{D}_W \mathbb{I}_n + 2\,\mathbb{D}_W \bigg) \cdot \nabla \operatorname{tr}\,\mathbb{D}_W \\ + \frac{1}{\operatorname{tr}\,\mathbb{D}_W} \bigg(\nabla \operatorname{tr}\,\mathbb{D}_W + 2\nabla \cdot \,\mathbb{D}_W \bigg) \bigg].$$

The values of the water diffusion tensor \mathbb{D}_W are known from DTI measurements.



A useful lemma

Lemma (Hillen, 2005)

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More general models and theory issues Consider the mean of velocity tensors $ar{f v}^{i_1\ldots i_k}:=\int_V v^{i_1}\ldots v^{i_k}dm v$

- If $k \in \mathbb{N}$ is odd, then $\overline{\mathbf{v}}^{i_1 \dots i_k} = \mathbf{0}, \ \forall \ i_1, \dots, i_k \in \{1, \dots, n\}.$
- If $k \in \mathbb{N}$ is even, then there is a constant $c_k > 0$ s.t.

$$\bar{\mathbf{v}}^{i_1\ldots i_k} = s^{k+n-1}c_k\Big(\sum_{\pi(i_1,\ldots,i_k)}\delta^{i_{j_1}i_{j_2}}\ldots\delta^{i_{j_k-1}i_{j_k}}\Big),$$

where the set of all pairs of indices out of (i_1, \ldots, i_k) is

 $\pi(i_1,\ldots,i_k):=\{((i_{j_1},i_{j_2}),\ldots,(i_{j_{k-1}},i_{j_k})) : \{j_1,\ldots,j_k\}=\{1,\ldots,k\}\}.$

The constants c_k are given by

$$c_0 = |\mathbb{S}^{n-1}|, \quad , c_2 = \frac{1}{n}|\mathbb{S}^{n-1}|, \quad , c_k = \frac{c_{k-2}}{k-2+n}, \text{ for } k \geq 4.$$



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More genera models and theory issues Setting $\mathbf{v} = s\theta$ with $\theta \in \mathbb{S}^{n-1}$ we get (let $\omega_n := |\mathbb{S}^{n-1}|$)

$$\mathbb{D}_{T}(\mathbf{x}) = \frac{1}{\lambda_{0}\omega} \int_{V} \mathbf{v} \otimes \mathbf{v} q d\mathbf{v}$$

$$= \frac{s^{2}}{\omega\lambda_{0}} \int_{\mathbb{S}^{n-1}} \theta \otimes \theta \frac{n}{\omega_{n} \operatorname{tr}\mathbb{D}_{W}(\mathbf{x})} \theta^{\dagger} \mathbb{D}_{W}(\mathbf{x}) \theta d\theta$$

$$= \frac{ns^{2}}{\omega\lambda_{0}\omega_{n} \operatorname{tr}\mathbb{D}_{W}} \int_{\mathbb{S}^{n-1}} \sum_{i,j}^{n} \theta^{i} \theta^{j} \sum_{k,l}^{n} \theta^{k} \theta^{l} D_{w}^{kl} d\theta$$

$$= \frac{ns^{2}}{\omega\lambda_{0}\omega_{n} \operatorname{tr}\mathbb{D}_{W}} \sum_{i,j,k,l}^{n} D_{w}^{kl} \int_{\mathbb{S}^{n-1}} \theta^{i} \theta^{j} \theta^{k} \theta^{l} d\theta$$



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By Lemma

$$\int_{\mathbb{S}^{n-1}} \theta^{i} \theta^{j} \theta^{k} \theta^{l} d\theta = c_{4} (\delta^{ij} \delta^{kl} + \delta^{ki} \delta^{lj} + \delta^{kj} \delta^{il})$$
$$= \frac{\omega_{n}}{n(n+2)} (\delta^{ij} \delta^{kl} + \delta^{ki} \delta^{lj} + \delta^{kj} \delta^{il})$$

thus

$$\mathbb{D}_{T}(\mathbf{x}) = \frac{ns^{2}}{\omega\lambda_{0}\omega_{n}\mathrm{tr}\mathbb{D}_{W}} \frac{\omega_{n}}{n(n+2)} \sum_{i,j,k,l}^{n} D_{w}^{kl} (\delta^{ij}\delta^{kl} + \delta^{ki}\delta^{lj} + \delta^{kj}\delta^{il})$$
$$= \frac{s^{2}}{\omega\lambda_{0}(n+2)\mathrm{tr}\mathbb{D}_{W}} \sum_{i,j}^{n} (\mathrm{tr}\mathbb{D}_{W} \ \delta^{ij} + D_{w}^{ij} + D_{w}^{ji}).$$

Assume \mathbb{D}_w is symmetric. This implies

$$\mathbb{D}_{T}(\mathbf{x}) = \frac{s^{2}}{\omega \lambda_{0}(n+2)} \left(\mathbb{I}_{n} + 2 \frac{\mathbb{D}_{W}}{\operatorname{tr} \mathbb{D}_{W}} \right).$$



Simulation results (with peanut)

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day 0, FA



day 0, Q



day 200, FA



day 200, Q



day 400, FA



day 400, Q



day 600, FA



day 600, Q



Simulation results in 3D (with estmated *Q* and peanut)

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day 200





day 400





day 600





Comparison with a pure macroscopic model

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$$\partial_t M_0 - \nabla \cdot (\mathbb{D}_T(\mathbf{x}) \nabla M_0) = Q(\mathbf{x}) \mu(M_0) M_0.$$



day 200, Q



day 400, Q



day 600, Q



Alternative proliferation modeling: Go-or-grow

Moving cancer cells:

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$$\partial_t p + \nabla_{\mathbf{x}} \cdot (\mathbf{v}p) + \partial_y \left(G(y, Q)p \right) = \mathcal{L}[\lambda]p - a(\mathbf{x})p + \frac{bq}{\omega}r - \ell(N)p$$

.



Alternative proliferation modeling: Go-or-grow

Moving cancer cells:

 $\partial_t p + \nabla_{\mathbf{x}} \cdot (\mathbf{v}p) + \partial_y \left(G(y, Q)p \right) = \mathcal{L}[\lambda]p - a(\mathbf{x})p + \frac{bq}{\omega}r - \ell(N)p$

• Non-moving (proliferating) cancer cells:

$$\partial_t r = a(\mathbf{x}) \int_V p d\mathbf{v} - br + g(N)r - \ell(N)r.$$

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Alternative proliferation modeling: Go-or-grow

Moving cancer cells:

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• Non-moving (proliferating) cancer cells:

$$\partial_t r = a(\mathbf{x}) \int_V p d\mathbf{v} - br + g(N)r - \ell(N)r.$$

 $\mathcal{L}[\lambda]p := -\lambda(y)p + \lambda(y) \frac{q(\mathbf{x}, \hat{\mathbf{v}})}{\omega} \int_V p(\mathbf{v}') d\mathbf{v}'$ (turning operator)

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Alternative proliferation modeling: Go-or-grow

Moving cancer cells:

 $\partial_t p + \nabla_{\mathbf{x}} \cdot (\mathbf{v}p) + \partial_y \left(G(y, Q)p \right) = \mathcal{L}[\lambda]p - a(\mathbf{x})p + \frac{bq}{\omega}r - \ell(N)p$

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Proliferation via cell-tissue interactions

Alternative proliferation modeling: Go-or-grow Therapy

More general models and theory issues $\mathcal{L}[\lambda]p := -\lambda(y)p + \lambda(y)\frac{q(\mathbf{x}, \hat{\mathbf{v}})}{\omega} \int_{V} p(\mathbf{v}')d\mathbf{v}' \quad (\text{turning operator})$ • Subcellular (receptor) dynamics:

$$\frac{d}{dt}y(t)=G(y(t),Q),$$



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Alternative proliferation modeling: Go-or-grow

Moving cancer cells:

 $\partial_t p + \nabla_{\mathbf{x}} \cdot (\mathbf{v}p) + \partial_y \left(G(y, Q)p \right) = \mathcal{L}[\lambda]p - a(\mathbf{x})p + \frac{bq}{\omega}r - \ell(N)p$

• Non-moving (proliferating) cancer cells:

$$\partial_t r = a(\mathbf{x}) \int_V p d\mathbf{v} - br + g(N)r - \ell(N)r.$$

Proliferation v cell-tissue interactions Alternative

Alternative proliferation modeling: Go-or-grow Therapy

More general models and theory issues $\mathcal{L}[\lambda]p := -\lambda(y)p + \lambda(y)\frac{q(\mathbf{x},\hat{\mathbf{v}})}{\omega} \int_{V} p(\mathbf{v}')d\mathbf{v}' \quad \text{(turning operator)}$ • Subcellular (receptor) dynamics:

$$\frac{d}{dt}y(t)=G(y(t),Q),$$

• Total cell density (macroscopic):

$$N(t, \mathbf{x}) = \int_V \int_Y p(t, \mathbf{x}, \mathbf{v}, y) dy \ d\mathbf{v} + \int_Y r(t, \mathbf{x}, y) dy$$



Deducing the effective macroscale equations

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Assumptions:

- system is close to steady-state;
- time scale on which birth and death events occur is much slower than the (biased) random walk process.

Wanted moments:

$$M(t, \mathbf{x}) := \iint_{V \times Z} p(t, \mathbf{x}, \mathbf{v}, z) dz d\mathbf{v}$$
$$w(t, \mathbf{x}, \mathbf{v}) := \int_{Z} r(t, \mathbf{x}, \mathbf{v}, z) dz$$

to recover the macroscopic cell density $N(t, \mathbf{x})$.



Deducing the effective macroscale equations

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$$w(t, \mathbf{x}, \mathbf{v}) := \int_{Z} r(t, \mathbf{x}, \mathbf{v}, z) dz$$

to recover the macroscopic cell density $N(t, \mathbf{x})$. Parabolic scaling again: $\hat{t} = \varepsilon^2 t$, $\hat{\mathbf{x}} = \varepsilon \mathbf{x}$.

$$g(N) \to \varepsilon^2 \hat{g}(\hat{N})$$

 $\ell(N) \to \varepsilon^2 \hat{\ell}(\hat{N}).$



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More general models and theory issues

We obtain

$$\partial_{t} N_{0} - \nabla \cdot \left(\frac{1}{\lambda_{0} + a(\mathbf{x})} \nabla \cdot \left(\frac{b}{a(\mathbf{x}) + b} \mathbb{D}_{T}(\mathbf{x}) N_{0}\right)\right) \\ + \nabla \cdot \left(\frac{\lambda_{1}}{\lambda_{0} + a(\mathbf{x})} \gamma(\mathbf{x}) f'(Q) \frac{b}{a(\mathbf{x}) + b} \mathbb{D}_{T}(\mathbf{x}) \cdot \nabla Q \ N_{0}\right) \\ = \frac{a(\mathbf{x})}{a(\mathbf{x}) + b} g(N_{0}) N_{0} - N_{0} \ell(N_{0}),$$



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We obtain

$$\begin{split} \partial_t N_0 &- \nabla \cdot \left(\frac{1}{\lambda_0 + a(\mathbf{x})} \nabla \cdot \left(\frac{b}{a(\mathbf{x}) + b} \mathbb{D}_T(\mathbf{x}) N_0 \right) \right) \\ &+ \nabla \cdot \left(\frac{\lambda_1}{\lambda_0 + a(\mathbf{x})} \gamma(\mathbf{x}) f'(Q) \frac{b}{a(\mathbf{x}) + b} \mathbb{D}_T(\mathbf{x}) \cdot \nabla Q \ N_0 \right) \\ &= \frac{a(\mathbf{x})}{a(\mathbf{x}) + b} g(N_0) N_0 - N_0 \ell(N_0), \end{split}$$

with the tumor diffusion tensor $\mathbb{D}_T(\mathbf{x}) = \frac{1}{\omega} \int_V \mathbf{v} \mathbf{v}^t q(\hat{\mathbf{v}}) d\mathbf{v}$.



Effective equations on the macroscale

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More general models and theory issues With the logistic growth choice $g(N_0) = c_g$, $\ell(N_0) = c_\ell N_0$, where $N_0 = \frac{a+b}{b}M_0$, we get

 $\partial_t N_0 - c_D(\mathbf{x}) \nabla \nabla \left(\mathbb{D}_T(\mathbf{x}) N_0 \right) - \lambda_1 c_D(\mathbf{x}) \nabla \left(\mathbf{u}(\mathbf{x}) N_0 \right) \\ = \frac{a}{a+b} c_g N_0 - c_\ell N_0^2,$

with $c_D(\mathbf{x}) = \frac{b}{(\lambda_0 + a(\mathbf{x}))(a(\mathbf{x}) + b)}$ and the drift velocity

 $\mathbf{u}(\mathbf{x}) = \gamma(\mathbf{x}) f'(Q(\mathbf{x})) \mathbb{D}_T(\mathbf{x}) \nabla Q,$

where $\gamma(\mathbf{x}) = (k^+Q + k^- + \lambda_0 + a)^{-1}$ and $f(Q(\mathbf{x})) = \frac{k^+Q(\mathbf{x})R_0}{k^+Q(\mathbf{x})+k^-}$.



Simulation results, full multiscale model

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Simulation results, model without proliferation

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Alternative proliferation modeling: Go-or-grow Therapy





Comparison between the two types of proliferation (time: 1 year, estimated Q, ODF)

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Proliferation vi cell-tissue interactions

Alternative proliferation modeling: Go-or-grow Therapy





Micro-meso model via go-or-grow, with therapy

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Proliferation via cell-tissue interactions

Alternative proliferation modeling: Go-or-grow Therapy

- Chemotherapy: inhibition of receptor binding (by peptidomimetics), with dosis *d_c*
- Radiotherapy: cell killing by ionizing radiation, with dosis d_r



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More general models and theory issues

- Chemotherapy: inhibition of receptor binding (by peptidomimetics), with dosis *d_c*
- Radiotherapy: cell killing by ionizing radiation, with dosis d_r

 $\partial_t p + \nabla_{\mathbf{x}} \cdot (\mathbf{v}p) + \partial_y \left(G(y, Q, d_c, d_r)p \right)$ = $\mathcal{L}[\lambda(y)]p - \mathbf{a}(\mathbf{x}, d_c)p + \mathbf{b}(\mathbf{x}, d_c) \frac{q(\hat{\mathbf{v}})}{\omega}r - L_1(N, \alpha_1, d_r)p$



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Proliferation via cell-tissue interactions

Alternative proliferation modeling: Go-or-grow Therapy

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$$\partial_t p + \nabla_{\mathbf{x}} \cdot (\mathbf{v}p) + \partial_y \left(G(y, Q, d_c, d_r)p \right)$$

= $\mathcal{L}[\lambda(y)]p - \mathbf{a}(\mathbf{x}, d_c)p + \mathbf{b}(\mathbf{x}, d_c) \frac{q(\hat{\mathbf{v}})}{\omega}r - L_1(N, \alpha_1, d_r)p$

$$\partial_t r = a(\mathbf{x}, d_c) \int_V p(\mathbf{v}) d\mathbf{v} - b(\mathbf{x}, d_c) r + g(N, d_c) r - L_2(N, \alpha_2, d_r) r$$



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Proliferation via cell-tissue interactions

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= $\mathcal{L}[\lambda(y)]p - \mathbf{a}(\mathbf{x}, d_c)p + \mathbf{b}(\mathbf{x}, d_c) \frac{q(\hat{\mathbf{v}})}{\omega}r - L_1(N, \alpha_1, d_r)p$

$$\partial_t r = a(\mathbf{x}, d_c) \int_V p(\mathbf{v}) d\mathbf{v} - b(\mathbf{x}, d_c) r + g(N, d_c) r - L_2(N, \alpha_2, d_r) r$$

with $L_l(N, \alpha_l, d_r) := \ell_l(N) + R_l(\alpha_l, d_r)$ $(l = 1, 2).$



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Alternative proliferation modeling: Go-or-grow Therapy

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$$\partial_t p + \nabla_{\mathbf{x}} \cdot (\mathbf{v}p) + \partial_y \left(G(y, Q, d_c, d_r)p \right)$$

= $\mathcal{L}[\lambda(y)]p - \mathbf{a}(\mathbf{x}, d_c)p + \mathbf{b}(\mathbf{x}, d_c) \frac{q(\hat{\mathbf{v}})}{\omega}r - L_1(N, \alpha_1, d_r)p$

$$\partial_t r = a(\mathbf{x}, d_c) \int_V p(\mathbf{v}) d\mathbf{v} - b(\mathbf{x}, d_c) r + g(N, d_c) r - L_2(N, \alpha_2, d_r) r$$

with $L_l(N, \alpha_l, d_r) := \ell_l(N) + R_l(\alpha_l, d_r)$ $(l = 1, 2)$.
 $\dot{y} = G(y, Q, d_c, d_r) = k^+ (d_c)(R_0 - y)Q S(\alpha_3, d_r) - k^- (d_c)y$.



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More genera models and theory issues $R_{j}(\alpha_{j}, d_{r}) = \sum_{i=1}^{\nu} (1 - S(\alpha_{j}, d_{r}))\eta_{\delta}(t - t_{i}), \qquad t_{i} \in \text{radiotherapy},$ supp $\eta_{\delta} \subset (-\delta, \delta), \ \delta << 1, \ j = 1, 2, 3.$ Survival fractions (LQ model): $S(\alpha_{j}, d_{r}) = \exp(-\alpha_{j}d_{r} - \beta_{j}d_{r}^{2}).$ For ν fractions, each of dosis \hat{d}_{r} :

 $S(\alpha_j, d_r) = \exp(-\nu(\alpha_j \hat{d}_r + \beta_j \hat{d}_r^2)) = \exp(-\alpha_j d_r (1 + \hat{d}_r/(\alpha_j/\beta_j))).$

- α_j represents lethal lesions produced by a single radiation track (α_jd_r, cell kill per Gy)
- β_j characterizes lethal lesions produced by two radiation tracks ($\beta_j d_r^2$, cell kill per Gy²)
- $\frac{\alpha_j}{\beta_j}$: radiation sensitivity, correlates with cell cycle length



Effective equation on macroscale

Multiscale models for glioma invasion

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More general models and theory issues

Remember $N(t, \mathbf{x}) = \int_V \int_Y p(t, \mathbf{x}, \mathbf{v}, y) dy \ d\mathbf{v} + \int_Y r(t, \mathbf{x}, y) dy.$



Effective equation on macroscale

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$$\partial_{t} N_{0} - \nabla \cdot \left(\frac{1}{\lambda_{0} + a} \nabla \cdot \left(\frac{b}{a + b} \mathbb{D}_{T}(\mathbf{x}) N_{0} \right) \right) \\ + \nabla \cdot \left(\frac{\lambda_{1} f'(Q)}{\gamma(\mathbf{x})} \frac{b}{(\lambda_{0} + a)(b + a)} \mathbb{D}_{T}(\mathbf{x}) \nabla Q N_{0} \right) \\ = \left(\left(g(N_{0}) - L_{2}(N_{0}) \right) \frac{a}{a + b} - L_{1}(N_{0}) \frac{b}{a + b} \right) N_{0},$$



Effective equation on macroscale

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Alternative proliferation modeling: Go-or-grow Therapy

More general models and theory issues Remember $N(t, \mathbf{x}) = \int_V \int_Y p(t, \mathbf{x}, \mathbf{v}, y) dy \ d\mathbf{v} + \int_Y r(t, \mathbf{x}, y) dy.$

$$\partial_{t} N_{0} - \nabla \cdot \left(\frac{1}{\lambda_{0} + a} \nabla \cdot \left(\frac{b}{a + b} \mathbb{D}_{T}(\mathbf{x}) N_{0} \right) \right) \\ + \nabla \cdot \left(\frac{\lambda_{1} f'(Q)}{\gamma(\mathbf{x})} \frac{b}{(\lambda_{0} + a)(b + a)} \mathbb{D}_{T}(\mathbf{x}) \nabla Q N_{0} \right) \\ = \left((g(N_{0}) - L_{2}(N_{0})) \frac{a}{a + b} - L_{1}(N_{0}) \frac{b}{a + b} \right) N_{0},$$

where $\gamma(\mathbf{x}) := \mathbf{k}^+ \mathbf{QS} + \mathbf{k}^- + \lambda_0 + \mathbf{a}$.



Therapy strategies

Multiscale models for glioma invasion

Christina Surulescu

Glioma invasion: Microscale, mesoscale, macroscale

Proliferation via cell-tissue interactions

Alternative proliferation modeling: Go-or-grow Therapy

- Strategy 1: resection (2 weeks after start), no further therapy.
- Strategy 2: resection (2 weeks after start), followed after 3 weeks by radiotherapy (weekends excluded) for 6 weeks.
- Strategy 3: resection (2 weeks after start), followed after 3 weeks by concurrent chemotherapy and radiotherapy (weekends excluded) for 6 weeks.


Results

Multiscale models for glioma invasion

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Proliferation vi cell-tissue interactions Alternative

proliferation modeling: Go-or-grow Therapy

More genera models and theory issues



start



before resection



after resection



Results

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end of therapy



end of therapy, scaled



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follow-up after two months



follow-up, scaled



Moving on...

Multiscale models for glioma invasion

Christina Surulescu

Glioma invasion: Microscale, mesoscale, macroscale

More general models and theory issues Goal: Multiscale model with tactic reorientations and tissue degradation.

- Involve on the mesolevel both hapto- and chemotaxis to describe cell reorientations → more complex subcellular dynamics.
- Model tissue degradation dependent on direction of cell motion (mesolevel description).



Subcellular dynamics revisited

Multiscale models for glioma invasion

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Glioma invasion: Microscale, mesoscale, macroscale

More general models and theory issues • Cells interact with the neighbouring tissue in order to move forward (contact guidance)



• Integrin receptor binding:



Subcellular dynamics revisited

Multiscale models for glioma invasion

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More general models and theory issues • Cells interact with the neighbouring tissue in order to move forward (contact guidance)



• to insoluble components Q

• Integrin receptor binding:

 $\frac{Q}{Q} + (R_0 - y_1 - y_2) \underset{k_1'}{\overset{k_1}{\rightleftharpoons}} RQ$

Notations: $y_1 := RQ$, $y_2 := RL$.



Subcellular dynamics revisited

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- Integrin receptor binding:
 - to insoluble components Q

$$egin{aligned} & \mathcal{Q} + (R_0 - y_1 - y_2) \stackrel{k_1}{\rightleftharpoons} RQ \ & \stackrel{k_1}{\mapsto} RQ, \ y_1 := RQ, \ y_2 := RL. \end{aligned}$$

• to soluble components L

$$L + (R_0 - y_1 - y_2) \stackrel{k_2}{\underset{k_2'}{\rightleftharpoons}} RL$$



Multiscale models for glioma invasion

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Glioma invasion: Microscale, mesoscale, macroscale

More general models and theory issues ODE for receptor dynamics ($j = 1, \ldots, N$, N: number of cells)

$$\frac{d\mathbf{y}^{(j)}}{dt} = \underbrace{\left(\begin{array}{c}k_1(R_0 - y_1 - y_2)Q(t, \mathbf{x}^{(j)}) - k_1'y_1\\k_2(R_0 - y_1 - y_2)L(t, \mathbf{x}^{(j)}) - k_2'y_2\end{array}\right)}_{\mathbf{G}(\mathbf{y}^{(j)}, Q(t, \mathbf{x}^{(j)}), L(t, \mathbf{x}^{(j)}))}$$

 R_0 total receptor concentration y_1, y_2 concentration of receptors bound to \overline{Q} , resp. L. In the absence of reorientations:

$$\frac{\partial p}{\partial t} + \underbrace{\mathbf{v} \cdot \nabla_{\mathbf{x}} p}_{\text{Transport with velocity } \mathbf{v}} + \underbrace{\nabla_{\mathbf{y}} \cdot (\mathbf{G}(\mathbf{y}, Q, L)p)}_{\text{Receptor dynamics}} = 0$$

Changes in orientation (and speed) have to be incorporated in the right-hand side.



Reorientation mechanisms - Haptotaxis

Multiscale models for glioma invasion

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Glioma invasion: Microscale, mesoscale, macroscale

More general models and theory issues

- Cells tend to align their movement direction to the direction of the fibers in the tissue.
- Haptotaxis:

$$\mathcal{H}(p,q)=\mathcal{H}_+(p,q)-\mathcal{H}_-(p,q).$$

• Gain term
$$\mathcal{H}_{+}(p,q) = \int_{V} \int_{\mathbb{S}^{n-1}} \eta_{h}(t,\mathbf{x},\mathbf{v}',\mathbf{y}) \psi(\mathbf{v};\mathbf{v}',\theta') p(\mathbf{v}') Q(\theta') d\mathbf{v}' d\theta'.$$

• Loss term

$$\mathcal{H}_{-}(p,q) = f(\mathbf{v}) \int_{V} \int_{\mathbb{S}^{n-1}} \eta_{h}(t, \mathbf{x}, \mathbf{v}, \mathbf{y}) \psi(\mathbf{v}'; \mathbf{v}, \theta') q(\theta') d\mathbf{v}' d\theta'.$$

 η_h is the rate for haptotactic reorientation $\psi(\mathbf{v}; \mathbf{v}', \theta')$ probability kernel for a reorientation $\mathbf{v}' \to \mathbf{v}$ after encounter with a fibre in direction θ' . E.g.,

 $\psi(\mathbf{v};\mathbf{v}',\theta') = |\hat{\mathbf{v}}'\cdot\theta'| \ \mathcal{K}_{H}^{(1)}(\mathbf{v},\theta') + (1-|\hat{\mathbf{v}}'\cdot\theta'|) \ \mathcal{K}_{H}^{(2)}(\mathbf{v},\mathbf{v}').$



Reorientation mechanisms - Chemotaxis

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Glioma invasion: Microscale, mesoscale, macroscale

More general models and theory issues • Chemotaxis: $C = C_+ - C_-$.

• Gain term
$$\mathcal{C}_{+}(p, L, \mathbf{y}) = \int_{V} \eta_{c}(t, \mathbf{x}, \mathbf{v}', \mathbf{y}) \mathcal{K}[L](\mathbf{v}, \mathbf{v}', \mathbf{y}) p(\mathbf{v}') d\mathbf{v}'.$$

Loss term

$$\mathcal{C}_{-}(\boldsymbol{p},\boldsymbol{L},\mathbf{y}) = \eta_{c}(t,\mathbf{x},\mathbf{v},\mathbf{y})\boldsymbol{p}(\mathbf{v}).$$

Turning kernel

$$\begin{split} & \mathcal{K}[L](\mathbf{v},\mathbf{v}',\mathbf{y}) = \alpha_1(\mathbf{y})\mathcal{K}(\mathbf{v},\mathbf{v}') + \alpha_2(\mathbf{y})\mathcal{K}(\mathbf{v},\nabla L).\\ & \alpha_1,\alpha_2: Y \to [0,1] \text{ with } \alpha_1(\mathbf{y}) + \alpha_2(\mathbf{y}) = 1, \ \forall \mathbf{y} \in Y. \end{split}$$



Full mesoscopic equation for cells

Multiscale models for glioma invasion

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Glioma invasion: Microscale, mesoscale, macroscale

More general models and theory issues



The macroscopic population density at time t and position \mathbf{x} is given by

$$M(t,\mathbf{x}) := \int_{Y} \int_{V} p(t,\mathbf{x},\mathbf{v},\mathbf{y}) d\mathbf{v} d\mathbf{y}$$



Tissue modification (Hillen 2005)

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Glioma invasion: Microscale, mesoscale, macroscale

More general models and theory issues

- Cell motion is both based on and can be impeded by ECM;
- Cells cut ECM-fibres using enzymes (proteolysis);
- Mean projection of movement direction on fiber orientation:

$$\Pi[p](t, \mathbf{x}, \theta) = \frac{1}{M(t, \mathbf{x})} \begin{cases} \int\limits_{Y} \int\limits_{V} |\theta \cdot \hat{\mathbf{v}}| p(t, \mathbf{x}, \mathbf{v}, \mathbf{y}) d\mathbf{v} d\mathbf{y}, & \text{undirected} \\ \int\limits_{Y} \int\limits_{V} \theta \cdot \hat{\mathbf{v}} p(t, \mathbf{x}, \mathbf{v}, \mathbf{y}) d\mathbf{v} d\mathbf{y}, & \text{directed} \end{cases}$$





Multiscale models for glioma invasion

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More general models and theory issues • Tissue equation based on mass-action kinetics:

$$\frac{\partial q}{\partial t} = \kappa(\Pi[p](t, \mathbf{x}, \theta) - 1)M(t, \mathbf{x})q(t, \mathbf{x}, \theta).$$

• Tissue degradation leads to production of a soluble ligand which then diffuses and degrades:

$$\frac{\partial L}{\partial t} = D_L \triangle L + \int_{\mathbb{S}^{n-1}} \kappa(1 - \Pi[p](t, \mathbf{x}, \theta)) M(t, \mathbf{x}) q(t, \mathbf{x}, \theta) d\theta - r_L h$$

 D_L : diffusion coefficient of L r_L : degradation rate of L κ : rate for proteolytic degradation of q.



Multiscale cell migration model

Multiscale models for glioma invasion

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Glioma invasion: Microscale, mesoscale, macroscale

More general models and theory issues

• Cells
$$p: [0, T] \times \mathbb{R}^n \times V \times Y \to \mathbb{R}$$

$$\frac{\partial p}{\partial t} + \mathbf{v} \cdot \nabla_{\mathbf{x}} p + \nabla_{\mathbf{y}} \cdot (\mathbf{G}(\mathbf{y}, Q, L)p) = \mathcal{H}(p, q) + \mathcal{C}(p, \nabla L)$$

with $p(0, \mathbf{x}, \mathbf{v}, \mathbf{y}) = p_0(\mathbf{x}, \mathbf{v}, \mathbf{y})$ and $\partial_n p = 0$ on ∂Y . • Tissue $q : [0, T] \times \mathbb{R}^n \times \mathbb{S}^{n-1} \to \mathbb{R}$

$$rac{\partial q}{\partial t} = \kappa (\Pi[p](t, \mathbf{x}, \theta) - 1) M(t, \mathbf{x}) q(t, \mathbf{x}, \theta)$$

with
$$q(0, \mathbf{x}, \theta) = q_0(\mathbf{x}, \theta)$$
.

• Soluble product of fibre cutting $L: [0, T] \times \mathbb{R}^n \to \mathbb{R}$

$$\frac{\partial L}{\partial t} = D_L \triangle L + \int_{\mathbb{S}^{n-1}} \kappa (1 - \Pi[p](t, \mathbf{x}, \theta)) M(t, \mathbf{x}) q(t, \mathbf{x}, \theta) d\theta - r_L L$$

with
$$L(0, \mathbf{x}) = L_0(\mathbf{x}) = 0$$



Global existence and uniqueness

Multiscale models for glioma invasion

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Glioma invasion: Microscale, mesoscale, macroscale

More general models and theory issues

Theorem (Kelkel & S. 2012)

There exists a global unique solution of the multiscale model in $L^1 \cap L^\infty$.

Theorem (Lorenz & S. 2014)

There exists a global unique solution of a more general multiscale model (allowing e.g., for nonlocal cell-tissue interactions) in L^2 .

Nonlocal cell-tissue interactions:

$$\frac{\partial q}{\partial t} = \kappa (\Pi[p](t, \mathbf{x}, \theta) - 1) M(t, \mathbf{x}) \mathcal{K} \star \frac{q(t, \mathbf{x}, \theta)}{1 + \gamma_Q |q(t, \mathbf{x}, \theta)|},$$

 $\mathcal{K} \star \phi$: convolution of a given spatial kernel \mathcal{K} with a function ϕ .



Conclusions

Multiscale models for glioma invasion

Christina Surulescu

Glioma invasion: Microscale, mesoscale, macroscale

More general models and theory issues

Multiscale models:

- allow testing the influence of many factors;
- are more difficult to handle numerically and analytically:
 - high dimensionality;
 - different scales both w.r.t. space and time;
 - highly nonlinear coupling;



Conclusions

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More general models and theory issues

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- allow testing the influence of many factors;
- are more difficult to handle numerically and analytically:
 - high dimensionality;
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 - highly nonlinear coupling;

Micro-meso-macro models:

- Other cells (no tissue): Erban & Othmer (Multiscale Model. Simul. 2005); Xue & Othmer (SIAP 2009).
- Tumor cells (moving in tissue networks): Bellomo et al. (M3AS 2010); Kelkel & S. (MBE 2011); Kelkel & S. (M3AS 2012); Lorenz & S. (M3AS 2014); Engwer, Hillen, Knappitsch, S. (JMB 2015); Engwer, Hunt & S. (IMA Math. Med. Biol. 2015); Engwer, Knappitsch & S. (MBE 2015); Hunt & S. (2015).



Outlook

Multiscale models for glioma invasion

Christina Surulescu

Glioma invasion: Microscale, mesoscale, macroscale

More general models and theory issues

- Goal: predict better CTVs and PTVs to allow for patient-specific treatment planning.
- tissue degradation modeled on mesoscale → chemotaxis equations on macroscale?
- lymph and blood angiogenesis;
- cell-cell interactions → effects on proliferation & invasion; can we recover macroscopic adhesion models?
- effects of hypoxia, acid-mediated invasion, tumor heterogeneity w.r.t. treatment response.