

**Problem 5: Hartree-Fock approximation****(3 points)**

The Hamilton operator of interacting electrons in a potential  $V(\vec{r})$  has the form

$$\hat{H} = \sum_{j=1}^N \left( \frac{\hat{p}_j^2}{2m} + V(\vec{r}_j) \right) + \frac{1}{2} \sum_{j \neq j'} \frac{e^2}{4\pi\epsilon_0} \frac{1}{|\vec{r}_j - \vec{r}_{j'}|} .$$

In the occupation number representation (basis:  $\psi_l(\vec{r}, \vec{s}) = \psi_l(\vec{x})$ ), the Hamilton operator is given by

$$\hat{H} = \sum_{i,l} A_{il} \hat{c}_i^+ \hat{c}_l + \sum_{i,j,l,m} \frac{1}{2} B_{ijlm} \hat{c}_i^+ \hat{c}_j^+ \hat{c}_l \hat{c}_m .$$

- a) Which relation exists between  $B_{ijlm}$  and  $B_{jiml}$ ?
- b) In the Hartree-Fock approach,  $\hat{H}$  is approximated by the effective Hamilton operator

$$\hat{H}_{\text{eff}} = \hat{H}_{\text{HF}} + W \quad \text{with} \quad \hat{H}_{\text{HF}} = \sum_{i,l} D_{il} \hat{c}_i^+ \hat{c}_l .$$

The term  $W$  does not contain any creation or annihilation operators. The approximation is realized by the following substitution of the product of four operators:

$$\begin{aligned} \hat{c}_i^+ \hat{c}_j^+ \hat{c}_l \hat{c}_m &\approx \hat{c}_i^+ c_m \langle \hat{c}_j^+ c_l \rangle_0 + \hat{c}_j^+ \hat{c}_l \langle \hat{c}_i^+ \hat{c}_m \rangle_0 \\ &- \langle \hat{c}_i^+ \hat{c}_m \rangle_0 \langle \hat{c}_j^+ \hat{c}_l \rangle_0 \\ &- \hat{c}_i^+ \hat{c}_l \langle \hat{c}_j^+ c_m \rangle_0 - \hat{c}_j^+ \hat{c}_m \langle \hat{c}_i^+ \hat{c}_l \rangle_0 \\ &+ \langle \hat{c}_i^+ \hat{c}_l \rangle_0 \langle \hat{c}_j^+ \hat{c}_m \rangle_0 . \end{aligned}$$

The expectation values  $\langle \rangle_0$  are calculated with the eigenfunctions of the effective Hamilton operator. Then:

$$\langle \hat{c}_i^+ c_m \rangle_0 = n_m \cdot \delta_{i,m} .$$

- i) Calculate  $D_{il}$  and  $W$ . Represent your result in terms of  $A_{il}$ ,  $B_{ijjl}$ ,  $B_{ijlj}$  and  $n_j$ .  
*Hint:* Use an appropriate substitution for the variables in the summation.
- ii) The coefficients  $D_{il}$  of the operator  $\hat{H}_{\text{HF}}$  can be written as

$$D_{il} = \int \psi_i^*(\vec{x}) \hat{O}(\vec{r}) \psi_l(\vec{x}) d^3x .$$

The functions  $\psi_l(\vec{x})$  are chosen to be eigenfunctions of  $\hat{O}$

$$\hat{O} \psi_l(\vec{x}) = \lambda_l \psi_l(\vec{x}) .$$

This leads to

$$D_{il} = \lambda_l \delta_{i,l} .$$

Compare this eigenvalue problem for  $\hat{O}$  with the Hartree-Fock equation in the position-spin representation (see last semester).

**Problem 6: Dielectric function of the electron gas****(4 points)**

The dielectric function of the three-dimensional electron gas has the form (Random phase approximation and  $T = 0$ )

$$\varepsilon(\vec{q}, \omega) = 1 - \frac{e^2}{\varepsilon_0 \Omega} \frac{2}{q^2} \sum_{\substack{\vec{k} \\ |\vec{k}| \leq k_F}} \left( \frac{1}{E(\vec{k}) - E(\vec{k} + \vec{q}) + \hbar\omega} + \frac{1}{E(\vec{k}) - E(\vec{k} + \vec{q}) - \hbar\omega} \right).$$

We consider the static limit  $\omega = 0$ .

- Calculate  $\varepsilon(\vec{q}, 0)$ . To this end, substitute the sum over  $\vec{k}$  by an integral.
- Consider the case  $q \ll 2k_F$  and take terms up to the order  $\frac{1}{q^2}$  into account. Calculate the screened potential in this case

$$\tilde{V}_{\text{eff}}(\vec{q}) = \frac{\tilde{V}_{\text{el}}(\vec{q})}{\varepsilon(\vec{q})} \quad \text{with} \quad \tilde{V}_{\text{el}}(\vec{q}) = -\frac{e^2}{\varepsilon_0 \Omega} \frac{1}{q^2}.$$

Use  $\tilde{V}_{\text{eff}}(\vec{q})$  to determine the screened potential  $V_{\text{eff}}(\vec{r})$  in real space.

- Discuss the behaviour of  $\varepsilon(\vec{q}, 0)$  in the limit  $q \rightarrow \infty$ ?
- Plot  $\varepsilon(\vec{q}, 0)$ .

*Hint:*

$$\int x \ln \left| \frac{ax + b}{ax - b} \right| = \frac{b}{a} x + \frac{1}{2} \left( x^2 - \frac{b^2}{a^2} \right) \ln \left| \frac{ax + b}{ax - b} \right|.$$

**Problem 7: Lorentz oscillator****(3 points)**

In a classical model for the screening in a solid, we assume that an external field

$$\vec{E}(t) = \vec{E}_0 \cdot e^{-i\omega t}$$

shifts the electrons in a solid by  $\vec{r}(t)$ . A “restoring force“  $-m\omega_0^2 \vec{r}(t)$  and the friction force  $-2m\gamma \dot{\vec{r}}(t)$  are acting in addition to the field  $\vec{E}$  on the electron. The displacement induces a dipol moment, which leads to a polarisation  $\vec{P} = -e \cdot n \vec{r}(t)$  with the electron density  $n$ .

- Solve Newton’s equation of motion for this system.
- Use

$$\varepsilon_0 \varepsilon(\omega) \vec{E}(\omega) = \varepsilon_0 \vec{E}(\omega) + \vec{P}(\omega)$$

to calculate the dielectric function.

Use the plasma frequency

$$\omega_p^2 = \frac{e^2}{\varepsilon_0} \frac{n}{m}$$

to represent your result.

- Decompose  $\varepsilon(\omega)$  into the real part  $\varepsilon_1(\omega)$  and the imaginary part  $\varepsilon_2(\omega)$ . Plot both functions.