Problem 8: Limit $\vec{q} \rightarrow 0$ of the dielectric function

The dielectric function $\varepsilon(\vec{q}, \omega)$ contains matrix elements of the form

$$I(\vec{q}\,) = \int \psi^*_{n',\vec{k}\,+\,\vec{q}}(\vec{r}\,) \,\mathrm{e}^{i\,\vec{q}\,r}\,\psi_{n,\vec{k}}\left(\vec{r}\,\right) d^3\,r \;.$$

They describe transitions between the bands n and n'. Consider the case of small wave vectors \vec{q} . The Bloch functions have the form

$$\psi_{n,\vec{k}}(\vec{r}) = e^{i\vec{k}\cdot\vec{r}} u_{n,k}(\vec{r})$$
 and $\psi_{n',\vec{k}+\vec{q}}(\vec{r}) = e^{i(\vec{k}+\vec{q})\cdot r} u_{n',k+q}(\vec{r})$.

- a) Use the lattice periodic functions $u_{n,\vec{k}}$ and $u_{n',\vec{k}+\vec{q}}$ to represent $I(\vec{q})$.
- b) The lattice periodic functions fulfill the following Schrödinger equations

$$\hat{H}\left(\vec{k}\,\right)u_{n,\vec{k}}\left(\vec{r}\,\right) \,=\, E_{n,\vec{k}}\,u_{n,\vec{k}}\left(\vec{r}\,\right)\,, \qquad \hat{H}\left(\vec{k}\,+\,\vec{q}\,\right)u_{n',\,k\,+\,\vec{q}}\left(r\right) \,=\, E_{n',\,\vec{k}\,+\,\vec{q}}\,u_{n',\,\vec{k}\,+\,q}\left(\vec{r}\,\right)\,.$$

Determine $\hat{H}(\vec{k})$ and $\hat{H}(\vec{k} + \vec{q})$.

c) $\hat{U} := \hat{H}(\vec{k} + \vec{q}) - \hat{H}(\vec{k})$ is a small perturbation for small wave vectors \vec{q} . Use first order perturbation theory to represent $u_{n',\vec{k}+\vec{q}}(\vec{r})$ by $u_{n,\vec{k}}(\vec{r})$. Employ this result to calculate $I(\vec{q})$. (*Hint*: the functions are orthonormal.) Use the matrix elements of the momentum operator

$$\vec{p}_{n',n}(\vec{k}\,) = \int u_{n',\vec{k}}^*(\vec{r}\,)\,\hat{\vec{p}}\,u_{n,\vec{k}}(\vec{r}\,)\,d^3\,r$$

to represent your final result.

Problem 9: Sum rule for dielectric function

(4 points)

A general property (which is called "sum rule") of the dielectric function is discussed in this problem.

- a) Consider a Hamilton operator \hat{H} with a complete orthonormal set of eigenfunctions $|\psi_{\alpha}\rangle$ and an additional operator \hat{A} .
 - i) Calculate the double commutator $[[\hat{H}, \hat{A}]_{-}, \hat{A}]_{-}$ and show that

$$\langle \psi_{\alpha} | [[\hat{H}, \hat{A}]_{-}, \hat{A}]_{-} | \psi_{\alpha'} \rangle = \sum_{\beta} \left(E_{\alpha} + E_{\alpha'} - 2 E_{\beta} \right) \left\langle \psi_{\alpha} | \hat{A} | \psi_{\beta} \right\rangle \left\langle \psi_{\beta} | \hat{A} | \psi_{\alpha'} \right\rangle \,.$$

ii) Use the general result of i) together with $\hat{A} = e^{-i\vec{q}\cdot\vec{r}}$ to calculate

$$S = \sum_{\beta} \left(E_{\alpha} - E_{\beta} \right) \left| \langle \psi_{\alpha} | \mathrm{e}^{-i \vec{q} \vec{r}} | \psi_{\beta} \rangle \right|^{2} \,.$$

(4 points)

b) The imaginary part of a dielectric function is given by

$$\begin{split} \varepsilon_{2}\left(\vec{q},\,\omega\right) \,&=\, \frac{\pi}{\Omega} \frac{\mathrm{e}^{2}}{\varepsilon_{0}\,q^{2}} \sum_{\delta,\,\vec{k},\,\vec{k}\,'\atop n,\,n'} f\left(E_{n,\,\vec{k}}\right) \, \left|\langle\psi_{n\,\vec{k}}|\mathrm{e}^{-i\,\vec{q}\,\vec{r}}|\psi_{n'\,\vec{k}\,\prime}\rangle\right|^{2} \\ &\cdot \left(\delta\left(E_{n\,\vec{k}}-E_{n'\,\vec{k}\,\prime}+\hbar\,\omega\right)-\delta\left(E_{n\,\vec{k}}-E_{n'\,\vec{k}\,\prime}-\hbar\,\omega\right)\right). \end{split}$$

A sum rule for $\varepsilon_2(\vec{q}, \omega)$ has the form

$$\int_{0}^{\infty} \omega \, \varepsilon_2 \left(\vec{q}, \, \omega \right) d \, \omega \, = \, \alpha \, \omega_p^2 \, .$$

Here, ω_p ist the plasma frequency. Use the result of a) to prove this sum rule and to calculate α .

c) The dielectric function in the Drude model has the form

$$\varepsilon(\omega) = \left(1 - \frac{\omega_p^2}{\omega^2 + i\gamma\omega}\right)$$
 with $\gamma > 0$ and $\gamma \to 0$.

 ω_p is the plasma frequency.

Calculate

a)
$$\int_{0}^{\infty} \omega \operatorname{Im}(\varepsilon(\omega)) d\omega$$
, b) $\int_{0}^{\infty} \operatorname{Im}\left(\frac{1}{\varepsilon(\omega)}\right) d\omega$

by explicit evaluation of the integrals in the limit $\gamma \rightarrow 0$.

Problem 10: Current density operator

(2 points)

The current density operator for a system of N particles with charge Q under the influence of a vector potential \vec{A} is defined in terms of the momentum and position operators $\hat{\vec{p}}_l = \frac{\hbar}{i} \vec{\nabla}_{\vec{r}_l}$ and $\hat{\vec{r}}_l$ of the *l*-th particle as follows

$$\vec{J}(\vec{r},t) = \hat{\vec{j}}(\vec{r}) + \hat{\vec{j}}_{dia}(\vec{r},t)$$

with a *paramagnetic*

$$\hat{\vec{j}}(\vec{r}) = \frac{1}{2} \frac{Q}{m} \sum_{l=1}^{N} \left(\hat{\vec{p}}_{l} \delta \left(\vec{r} - \hat{\vec{r}}_{l} \right) + \delta \left(\vec{r} - \hat{\vec{r}}_{l} \right) \hat{\vec{p}}_{l} \right)$$

and a *diamagnetic* part

$$\hat{\vec{j}}_{\text{dia}}\left(\vec{r},t\right) = -\frac{Q^2}{m} \sum_{l=1}^{N} \delta\left(\vec{r} - \hat{\vec{r}}_l\right) \vec{A}\left(\hat{\vec{r}}_l,t\right) \,.$$

i) Calculate $\hat{\vec{J}}(\vec{r}, t)$ in the occupation number representation for a basis of plane waves

$$\psi_{\vec{k},\sigma}\left(\vec{r}\right) = \frac{1}{\sqrt{\Omega}} e^{i\,\vec{k}\cdot\vec{r}} \chi_{\sigma} \,.$$

ii) Determine the Fourier transform of the operator resulting in i).